

Hankel and Toeplitz Determinants

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The most famous Hankel matrix is the Hilbert matrix

$$H_n = \frac{1}{i+j-1} \quad 1 \leq i, j \leq n$$

which has determinant equal to a ratio of Barnes G -function values:

$$\det(H_n) = \frac{\prod_{k=1}^{n-1} (k!)^4}{2^{n-1} \ell!} = \frac{G(n+1)^4}{G(2n+1)} \rightarrow 0$$

as $n \rightarrow \infty$. More precisely [1],

$$\frac{\det(H_n)}{4^{-n^2} (2\pi)^n n^{-1/4}} \rightarrow 2^{1/12} e^{1/4} A^{-3} = 0.6450024485\dots$$

where A denotes the Glaisher-Kinkelin constant [2]. Such Hankel determinants are important in random matrix theory and applications [3], but we shall forsake all this, giving instead only a few examples [4, 5, 6]. Another interesting fact is that $\det(H_n)$ is always the reciprocal of a positive integer [7].

The Hankel determinant of Euler numbers [8] is, in absolute value,

$$\begin{aligned} |E_{i+j}|_{0 \leq i, j \leq n-1} &= \prod_{k=1}^{n-1} (k!)^2 = G(n+1)^2 \\ &\sim \frac{e^{1/6}}{A^2} e^{-\frac{3}{2}n^2} (2\pi)^n n^{n^2 - \frac{1}{6}} \end{aligned}$$

as $n \rightarrow \infty$. The simplicity of this result contrasts with the following. The Hankel determinant of Bernoulli numbers [9] is, in absolute value,

$$\begin{aligned} |B_{i+j}|_{0 \leq i, j \leq n-1} &= \prod_{k=1}^{n-1} \frac{(k!)^6}{(2k)!(2k+1)!} \\ &= \frac{2^{\frac{1}{12}} e^{\frac{1}{4}}}{A^3} 4^{-n^2} (2\pi)^n \frac{G(n+1)^4}{G(n+1/2)G(n+3/2)} \\ &\sim \frac{2^{\frac{1}{12}} e^{\frac{5}{12}}}{A^5} 4^{-n^2} e^{-\frac{3}{2}n^2} (2\pi)^{2n} n^{n^2 - \frac{5}{12}} \end{aligned}$$

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as $n \rightarrow \infty$. We mention three formulas of Krattenthaler [10]:

$$\begin{aligned} \frac{B_{2i+2j+2}}{(2i+2j+2)!} &= 4^{-n^2} \prod_{k=1}^{2n-1} (2k+1)^{-2n+k}, \\ \frac{B_{2i+2j+4}}{(2i+2j+4)!} &= 4^{-n^2-n} 9^{-n} \prod_{k=1}^{2n-1} (2k+3)^{-2n+k}, \\ \frac{B_{2i+2j+6}}{(2i+2j+6)!} &= (n+1)(2n+3) 4^{-n^2-2n} \prod_{k=1}^{2n+1} (2k+1)^{-2n-2+k} \end{aligned}$$

which are always reciprocals of integers (unlike $|E_{i+j}|$ and $|B_{i+j}|$). The asymptotics of these three sequences remain open.

More difficult are determinants of Riemann zeta function values:

$$a_n^{(0)} = |\zeta(i+j)|_{1 \leq i, j \leq n}, \quad a_n^{(1)} = |\zeta(i+j+1)|_{1 \leq i, j \leq n}$$

which evidently satisfy

$$a_n^{(0)} \sim C \cdot \frac{\mu_{2n+1}}{e^{3/2}} \prod_{1 \leq i, j \leq n} \frac{1}{i-j}, \quad a_n^{(1)} \sim \frac{e^{9/8}}{\sqrt{6}} C \cdot \frac{\mu_{2n}}{e^{3/2}} \prod_{1 \leq i, j \leq n} \frac{1}{i-j}$$

thanks to numerical experiments by Zagier [11]. No closed-form expression for the constant $C = 0.351466738331\dots$ is known.

A famous Toeplitz matrix, called the alternating Hilbert matrix in [12], is

$$\tilde{H}_n = \prod_{1 \leq i, j \leq n} \frac{1}{i-j}$$

where we understand the diagonal elements to be 0. Schur [13] proved long ago that the maximum eigenvalue (in modulus) of both H_n and \tilde{H}_n is less than π and approaches π as $n \rightarrow \infty$. The determinant is, of course, the product of all eigenvalues. When n is odd, $\det(\hat{H}_n) = 0$. When n is even, a closed-form expression for $\det(\hat{H}_n)$ seems to be unavailable, despite the existence of a combinatorial approach [14]. Note that the ‘‘symbol’’ associated with \hat{H}_n is

$$\sum_{r=1}^{\infty} \frac{e^{ir\theta}}{-r} + \sum_{r=1}^{\infty} \frac{e^{-ir\theta}}{r} = i(\theta - \pi)$$

for $0 < \theta < 2\pi$, hence a theorem due to Grenander & Szegő [15] gives

$$\lim_{\substack{n \rightarrow \infty \\ n \text{ even}}} \frac{1}{n} \ln \det(\hat{H}_n) = \frac{1}{2\pi} \int_0^{2\pi} \ln [i(\theta - \pi)] d\theta = -1 + \ln(\pi) = 0.1447298858\dots$$

A refined estimate shown subsequently in [15], potentially governing the value of

$$\lim_{n \rightarrow \infty} \det(\hat{H}_n) \cdot \frac{\pi^n}{e^3},$$

has conditions that must be verified.

Consider finally another Toeplitz matrix

$$K_n = \frac{1}{1 + |i - j|} \quad 1 \leq i, j \leq n$$

for which little is known. The “symbol” here is

$$\sum_{r=0}^{\infty} \frac{e^{ir\theta}}{1+r} + \sum_{r=1}^{\infty} \frac{e^{-ir\theta}}{1+r} = -1 - e^{i\theta} \ln |1 - e^{-i\theta}| - e^{-i\theta} \ln |1 - e^{i\theta}|$$

for $0 < \theta < 2\pi$, hence

$$\begin{aligned} \lim_{n \rightarrow \infty} \frac{1}{n} \ln(\det(K_n)) &= \frac{1}{2\pi} \int_0^{2\pi} \ln | -1 - e^{i\theta} \ln |1 - e^{-i\theta}| - e^{-i\theta} \ln |1 - e^{i\theta}| | d\theta \\ &= -0.3100863233\dots \end{aligned}$$

An exact formula for this constant is desired; might, at least, the integral be simplified in some way?

0.1. Combinatorial Approach. Assume that n is even. Let S denote the set of all $(n/2)$ -tuples of ordered pairs:

$$(p_k, q_k)_{k=1}^{n/2}$$

of positive integers $p_k < q_k$ satisfying

$$\bigcup_{k=1}^{n/2} \{p_k, q_k\} = \{1, 2, \dots, n\}$$

and $p_1 < p_2 < \dots < p_{n/2}$. Note that the q s need not be in ascending order. Let us verify a formula in [14]:

$$\det(\hat{H}_n) = \prod_{(p_k, q_k)_{k=1}^{n/2} \in S} \prod_{k=1}^{n/2} \frac{1}{(q_k - p_k)^2}$$

for $n = 4$. Three such 2-tuples exist:

$$p_1 = 1 < p_2 = 2 < q_1 = 3 < q_2 = 4,$$

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yielding

$$\frac{1}{(3-1)^2(4-2)^2} + \frac{1}{(4-1)^2(3-2)^2} + \frac{1}{(2-1)^2(4-3)^2} = \frac{169}{144} = \det(\hat{H}_4).$$

The case $\det(\hat{H}_2) = 1$ is trivial; the case $\det(\hat{H}_6) = 6723649/4665600$ will require some effort. We wonder if a simple method for computing the size of S , as a function of n , can be found. An analogous approach for $\det(K_n)$ would also be good to see.

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